

# Symétrie $\mathcal{PT}$ , un concept unificateur pour la physique des dispositifs à gain

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Collaboration: A. Lupu<sup>2</sup>, and A. Degiron<sup>2</sup>

<sup>2</sup>IEF, Univ. Paris-Sud and CNRS, Orsay, France

H. Benisty, C. Yan, A. T. Lupu, and A. Degiron,  
*IEEE J. Lightwave Technol.*, vol. 30, pp. 2675-2683, 2012.

A. Lupu, H. Benisty and A. Degiron, *Optics Express* **21**, 21651 (2013)

H. Benisty and M. Besbes,  
*JOSA. B*, vol. 29, pp. 818-826, March 29 2012.

H. Benisty et al. *Opt. Express*, **19**, 18004, 2011

ENST 3 Oct 2013, Paris : Symétrie  $\mathcal{PT}$  Concept unificateur

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## OUTLINE

⇒  $\mathcal{PT}$ -symmetry, Hamiltonians, and Gain-Loss (with reminder)

⇒ Three flavours of  $\mathcal{PT}$ -symmetries

- With waveguides
- With resonators
- With gratings

⇒ Application with plasmonics (with IEF : A. Lupu, A. Degiron)

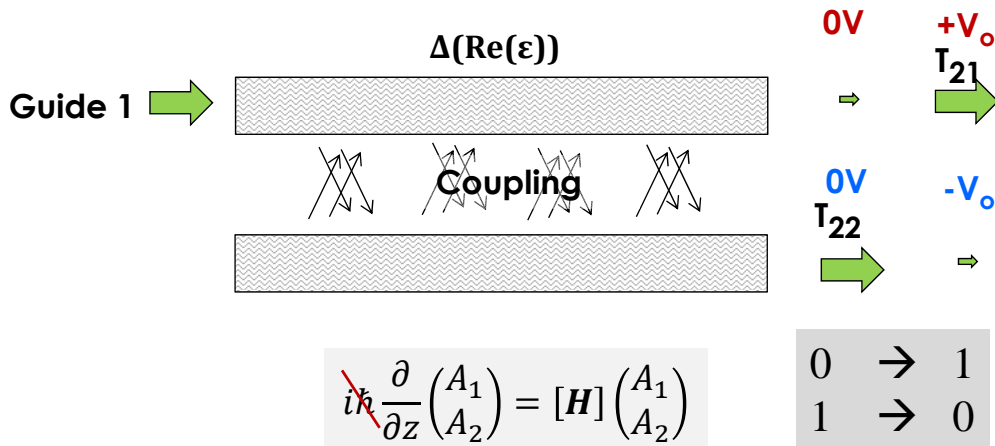
- Losses are now causing a singularity !
- Switches
- Real life ? (“healing” of smoothed singularity)

⇒  $\mathcal{PT}$ -symmetry and (many) current photonic concepts

- from coherent perfect absorbers to lasers

⇒ The brachistochrone problem (1696[!]-1990)

- Relevance for fast Quantum computation



$$\text{input} = \frac{1}{2} \begin{pmatrix} A \\ A \end{pmatrix} + \frac{1}{2} \begin{pmatrix} A \\ -A \end{pmatrix}$$

$$\text{output} = \frac{1}{2} \begin{pmatrix} A \\ A \end{pmatrix} e^{i\beta_+L} + \frac{1}{2} \begin{pmatrix} A \\ -A \end{pmatrix} e^{i\beta_-L}$$

controlled by  $(\beta_+ - \beta_-)L$

## Initial motivation generalized Quantum Mechanics

• Hermitian operator  $H$   $\rightarrow$  ♥ Real Eigenvalues

• Hermitian operator  $\leftarrow$  ♥ Real Eigenvalues

$$\begin{aligned} (\lambda - ig)(\lambda + ig) - \kappa^2 &= 0 \\ \lambda^2 &= \kappa^2 - g^2 \end{aligned}$$

• Counter-example (Bender 1998)

$$H = \begin{pmatrix} \hbar(\omega_1 + ig) & \kappa \\ \kappa & \hbar(\omega_2 - ig) \end{pmatrix}$$

•  $(\omega_1 = \omega_2 = 0)$  Real eigenvalues if :  $g < \kappa$

•  $\{1 \leftrightarrow 2\}$  &  $\{t \overset{\mathcal{T}}{\leftrightarrow} -t\} = 1$

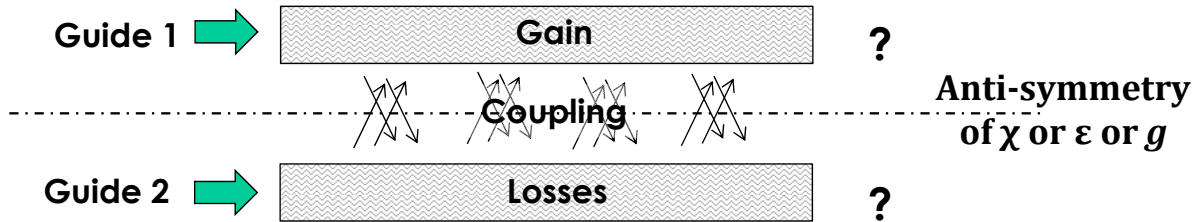
•  $\mathcal{PT}H = H$

$$i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

# Coupled waveguides !

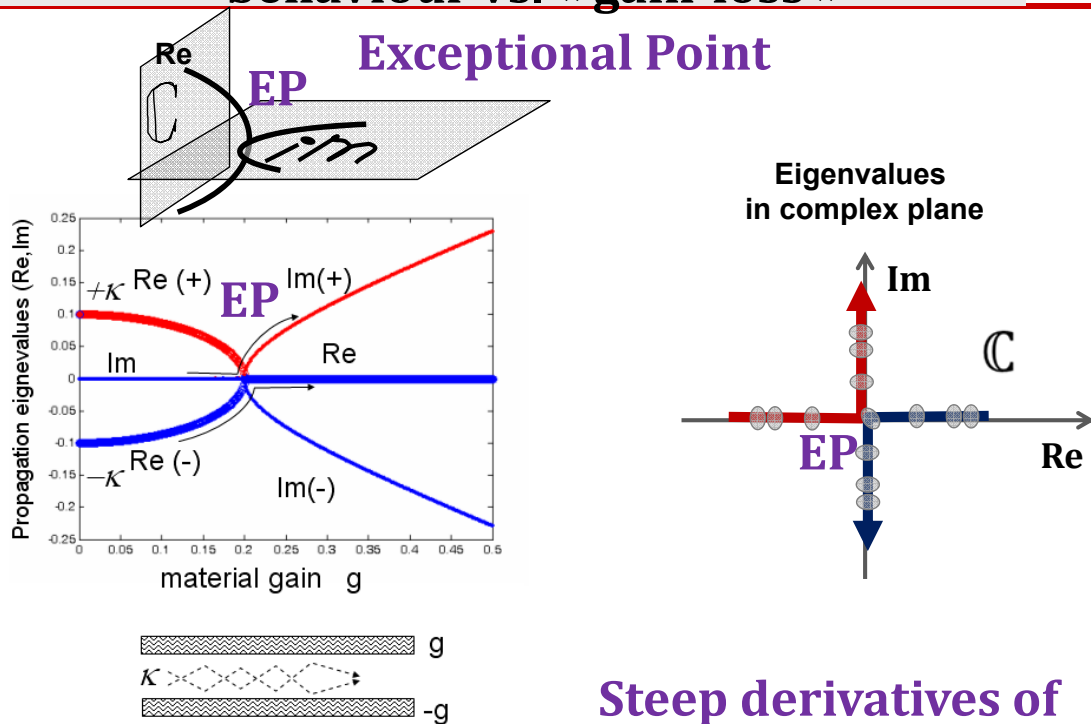
$$i\hbar \frac{\partial}{\partial z} \begin{pmatrix} A_1 \\ A_2 \end{pmatrix} = [H] \begin{pmatrix} A_1 \\ A_2 \end{pmatrix}$$

$$i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

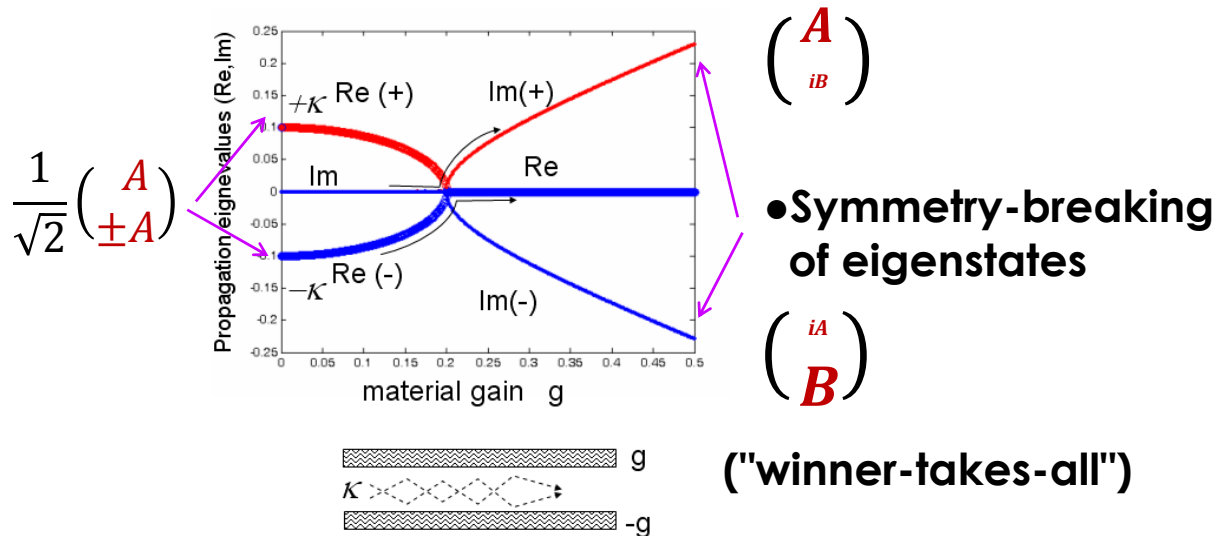


$$\begin{pmatrix} (\beta_1 + ig) & \kappa \\ \kappa & (\beta_2 - ig) \end{pmatrix}$$

# Canonical eigenvalue behaviour vs. « gain-loss »



## Symmetry breaking



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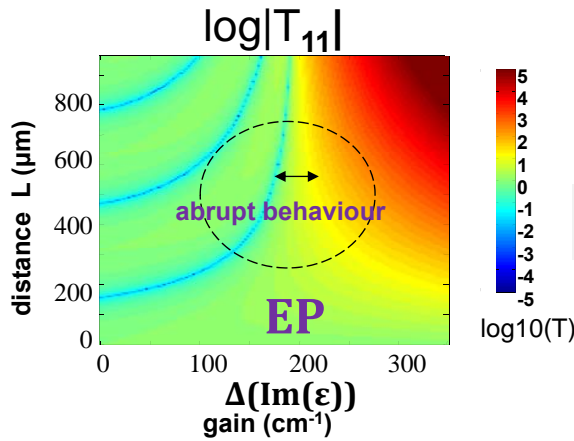
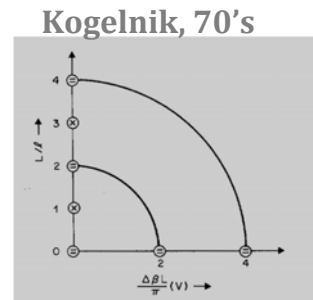
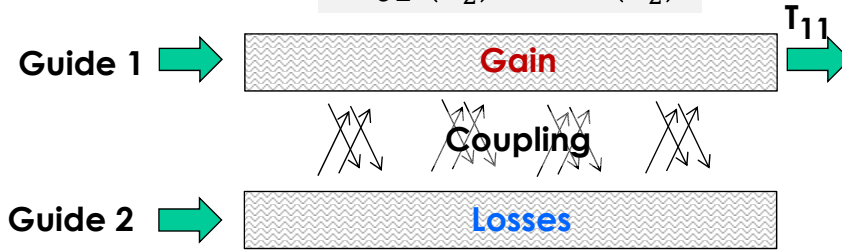
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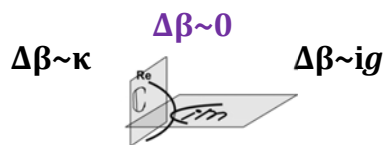
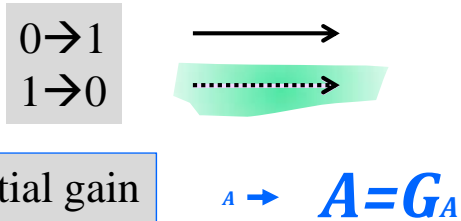
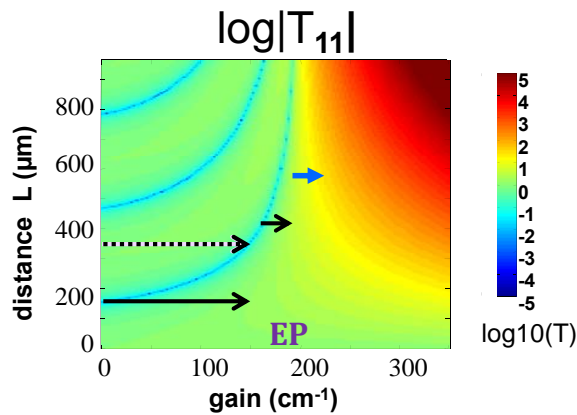
# PT-symmetry with two waveguides

$$i\hbar \frac{\partial}{\partial z} \begin{pmatrix} A_1 \\ A_2 \end{pmatrix} = [H] \begin{pmatrix} A_1 \\ A_2 \end{pmatrix}$$

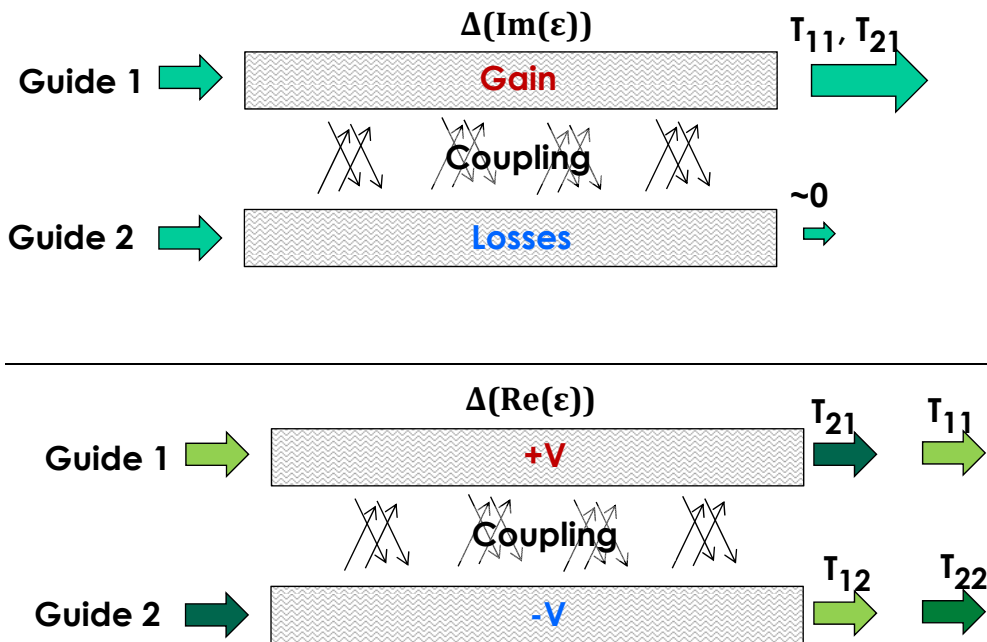


In Optics : ~ Ctyroky 1996 ! Klaiman 2008; --Guo 2009; --Rüter 2010,...

## Serves as switch or modulator



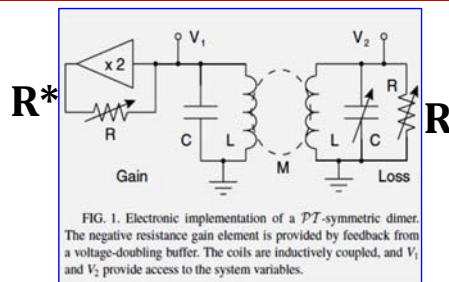
large  $\frac{\partial G}{\partial g}$



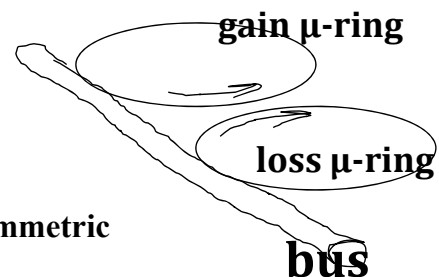
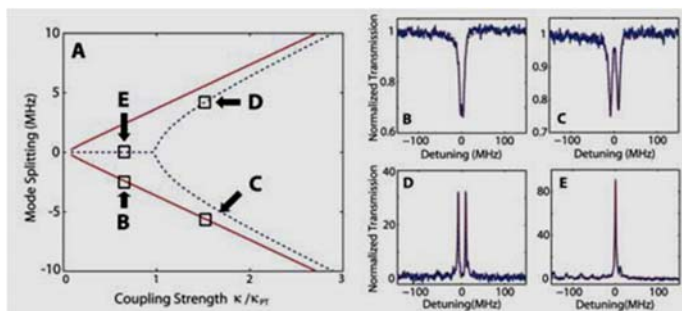
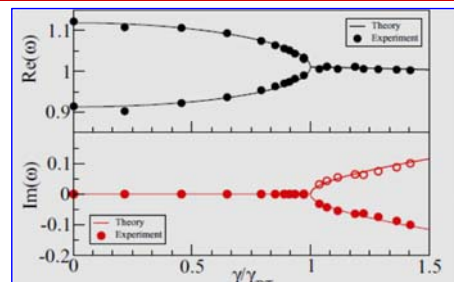
## PT-symmetry in two resonators

Easier !

$R^*LC$   
 $+RLC$



Phys. Rev. A, vol. 84,  
p. 040101, 2011.  
(Kottos, Christodoulos,...)



**Nonreciprocal light transmission in parity-time-symmetric whispering-gallery microcavities**

Bo Peng et al. Arxiv 2013 (coor author Lan Yang @ ese.wustl.edu)

PRL 108, 173901 (2012)

M. Lierzter,<sup>1,\*</sup> Li Ge,<sup>2</sup> A. Cerjan,<sup>3</sup> A. D. Stone,<sup>3</sup> H. E. Türeci,<sup>2,4</sup> and S. Rotter<sup>1,†</sup>

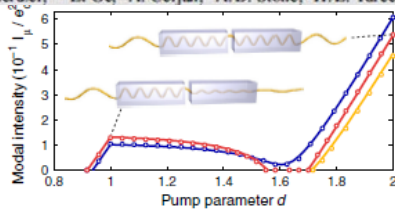
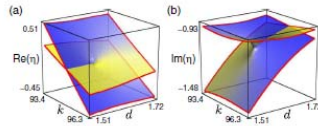
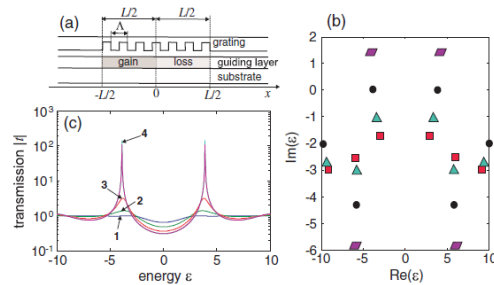


FIG. 1 (color online). Intensity output of a laser system consisting of two 1D coupled ridge lasers, each of length  $100 \mu\text{m}$  with an air gap of size  $10 \mu\text{m}$  and an (unpumped) index of refraction  $n = 3 + 0.13i$ . For  $0 < d < 1$ , the pump in the left ridge is linearly increased in the range  $0 < D < 1.2$ , and, for



Stefano Longhi  
PRL 105, 013903 (2010)

PHYSICAL REV



Non-Hermitian Dirac equation and its optical realization.—Let us consider the Dirac equation in one spatial



Won-Tien Tsang, one of the three inventors of the cleaved coupled-cavity laser, prepares

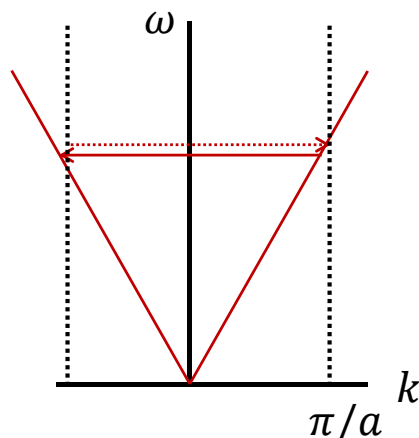
Bell Lab patent 1965  
Tsang 1984  
C3 laser  
(Coupled cavity laser)

## PT-symmetry with gratings

$$\Delta\epsilon(x) \sim \Delta\epsilon_1 [ (\exp iGx) + \exp -iGx) ]$$

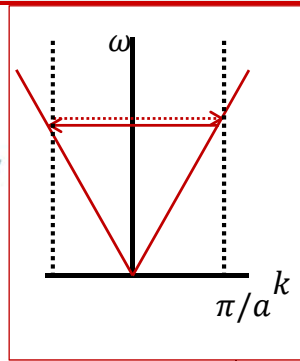
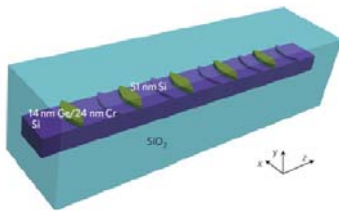
$$= \Delta\epsilon_1 [ \cos(kx) + i \sin(kx) ]$$

« Single sideband » grating (Fr:BLU)



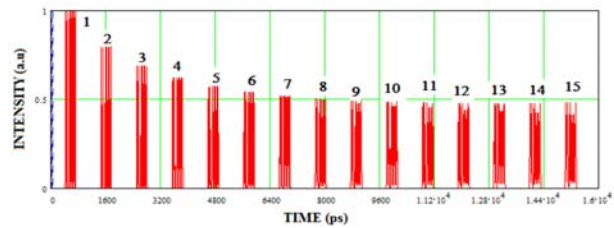
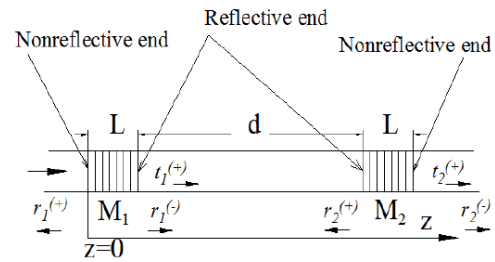


## (Fr:BLU)DBR



Kulishov et al.

22 April 2013 | Vol. 21, No. 8 | DOI:10.1364/OE.21.009473 | OPTICS EXPRESS 9473

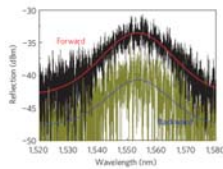


nature materials

LETTERS

Experimental demonstration of a unidirectional reflectionless parity-time metamaterial at optical frequencies

$$\Delta\epsilon(x) \sim \Delta\epsilon_1 [\cos(kx) + i(-1 + \sin(kx))]$$



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- Real life ? (“healing” of smoothed singularity)

⇒ **PT-symmetry and (many) current photonic concepts**

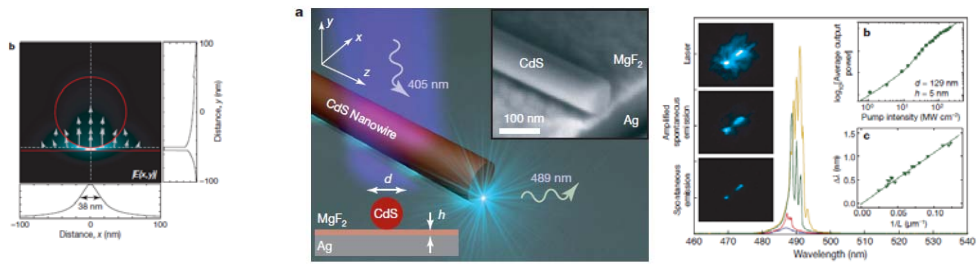
- from coherent perfect absorbers to lasers

⇒ **The brachistochrone problem (1696[!]-1990)**

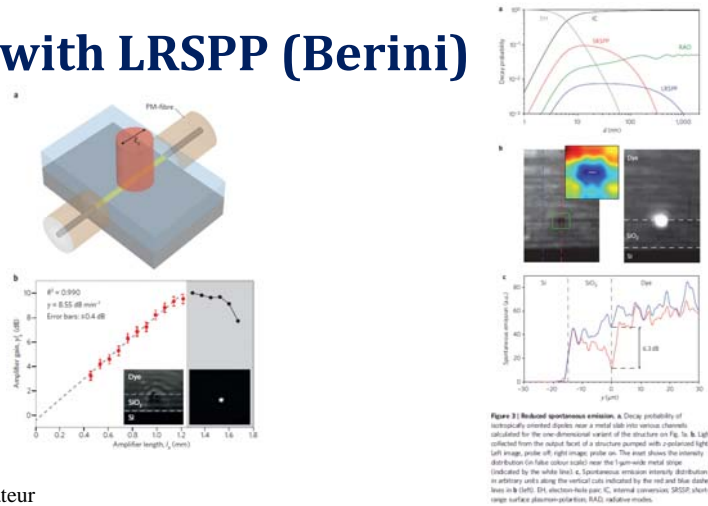
- Relevance for fast Quantum computation



## • SPASER (Stockman, Oulton with nanorods, ...)



## • Optical Amplifiers with LRSPP (Berini)



# PT-symmetry with relaxed gain-loss balance

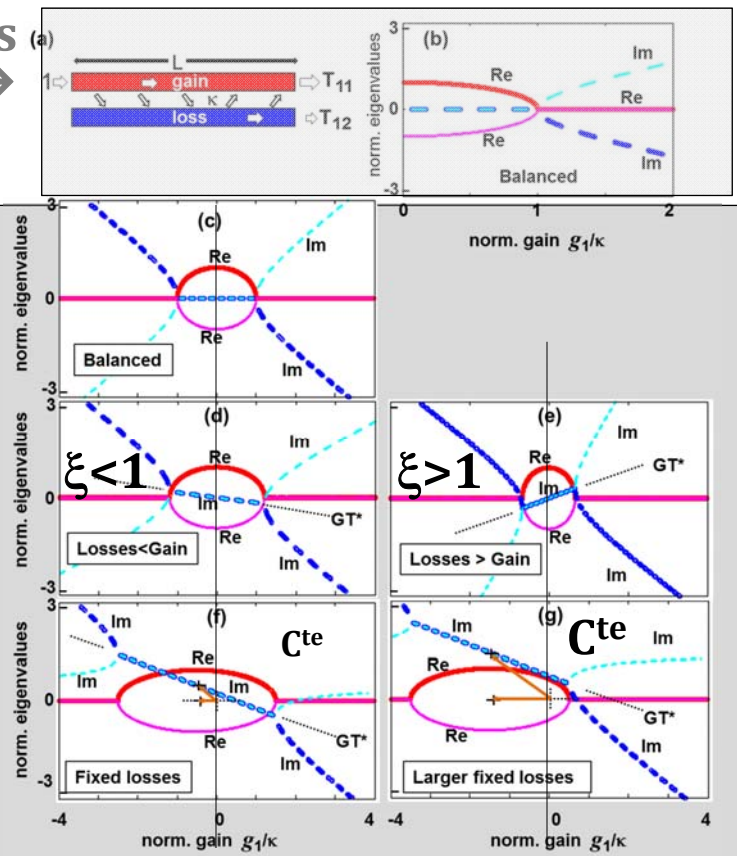
previous representation →

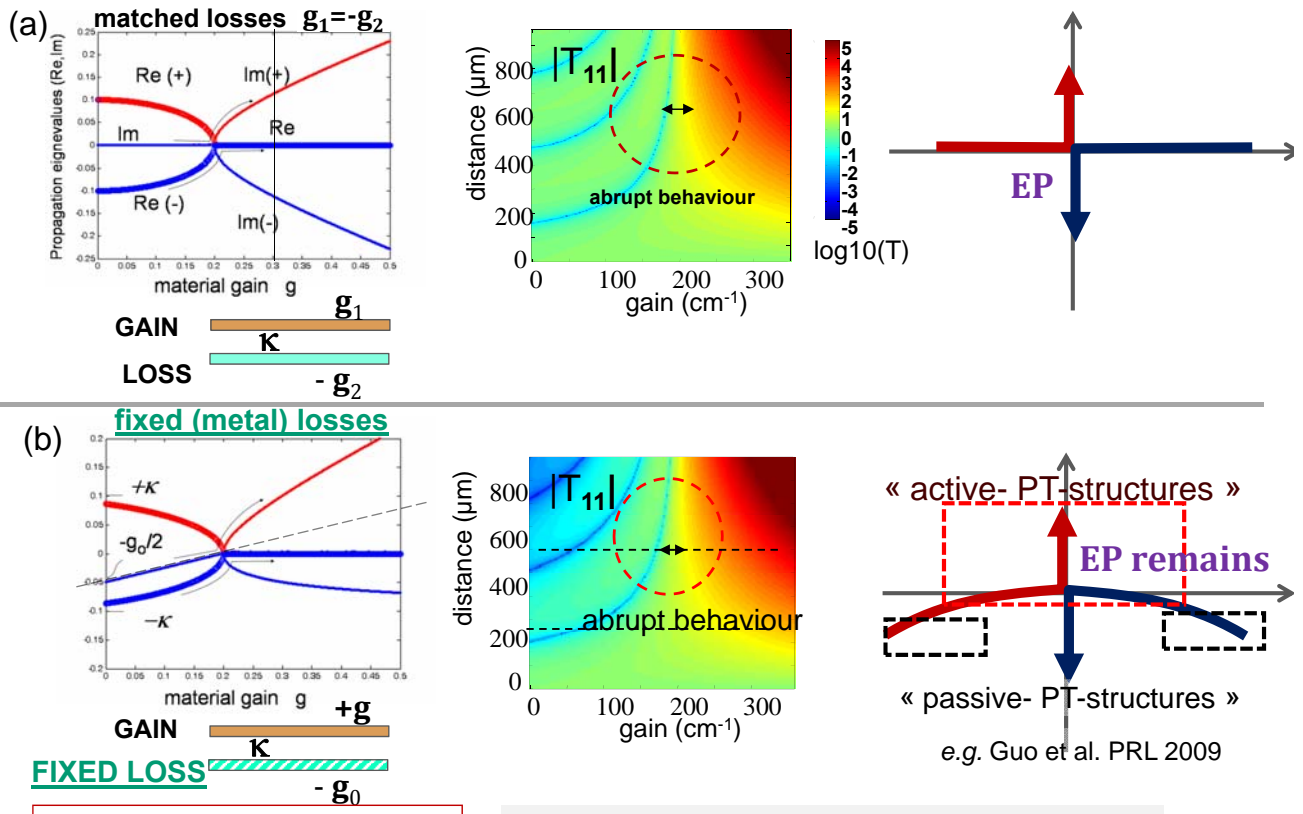
**broad view**

$$|\text{loss}| = \text{gain}$$

$$|\text{loss}| = \xi \times \text{gain}$$

$$|\text{loss}| = C^{\text{te}}$$

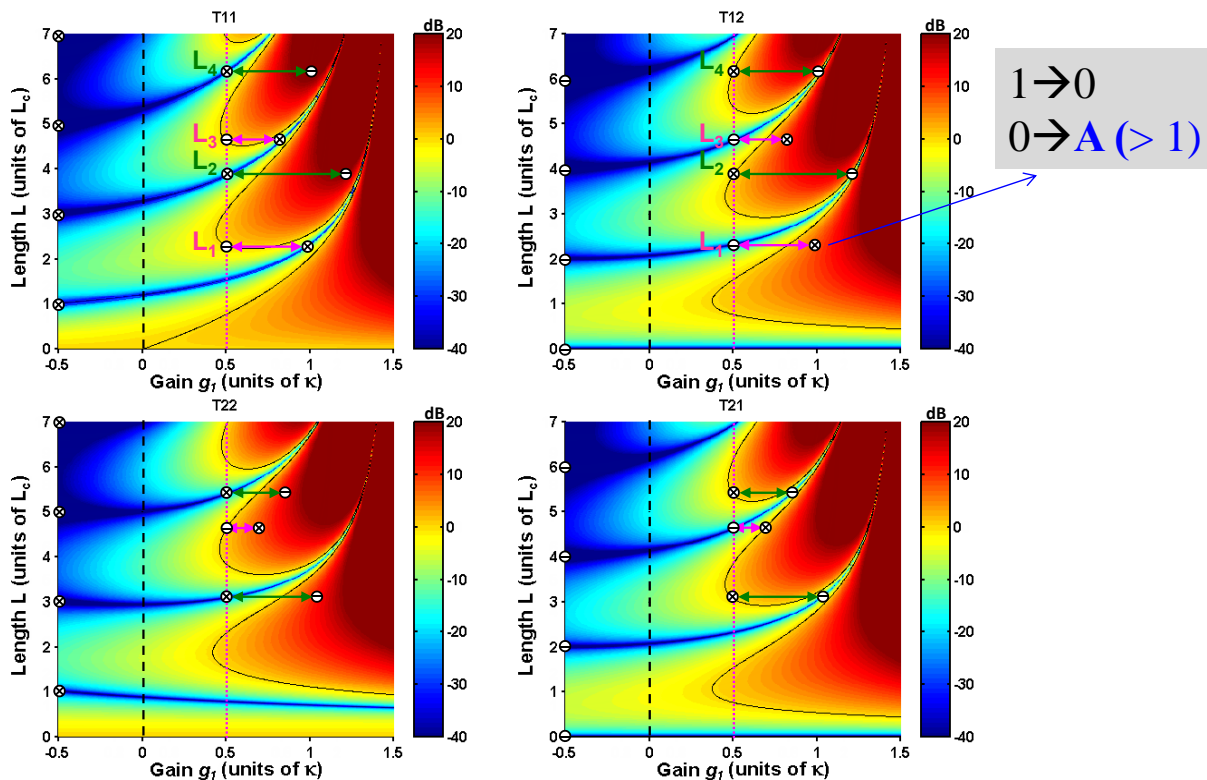




$g_0$  and  $\kappa$  should be "matched"

H. Benisty et al. *Opt. Express*, **19**, 18004, 2011

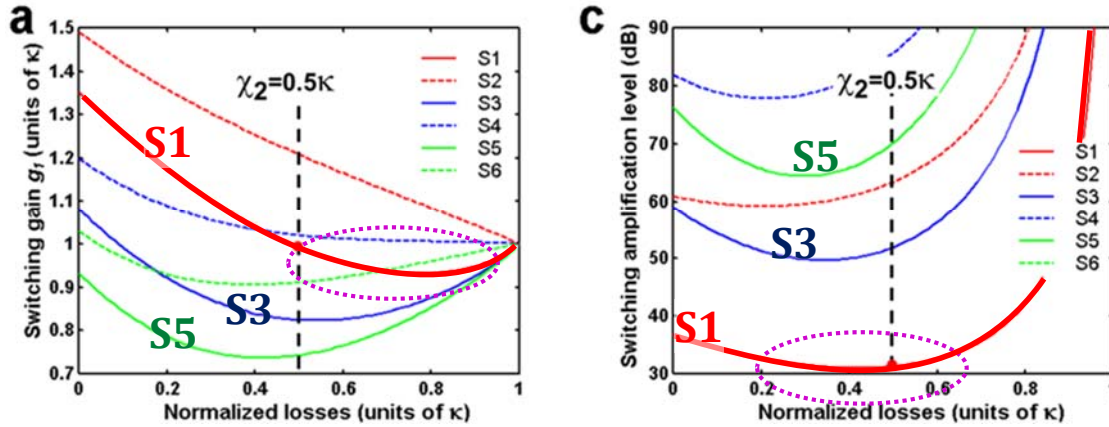
## Whole coupler picture



## Switching using PT symmetry in plasmonic systems: positive role of the losses

Anatole Lupu,<sup>1,2,\*</sup> Henri Benisty,<sup>3</sup> and Aloyse Degiron<sup>1,2</sup>

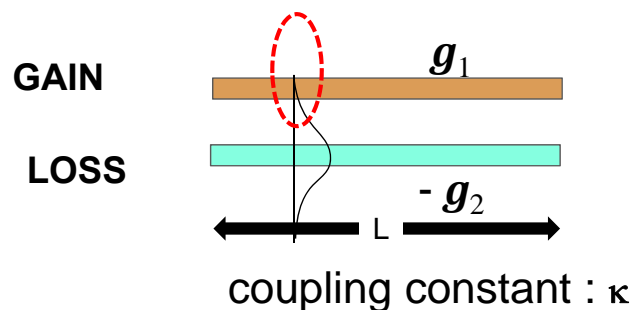
9 September 2013 | Vol. 21, No. 18 | DOI:10.1364/OE.21.021651 | OPTICS EXPRESS 21651



Existence of optimal losses, minimizing requirement for **gain & amplification** vs.  $\kappa$

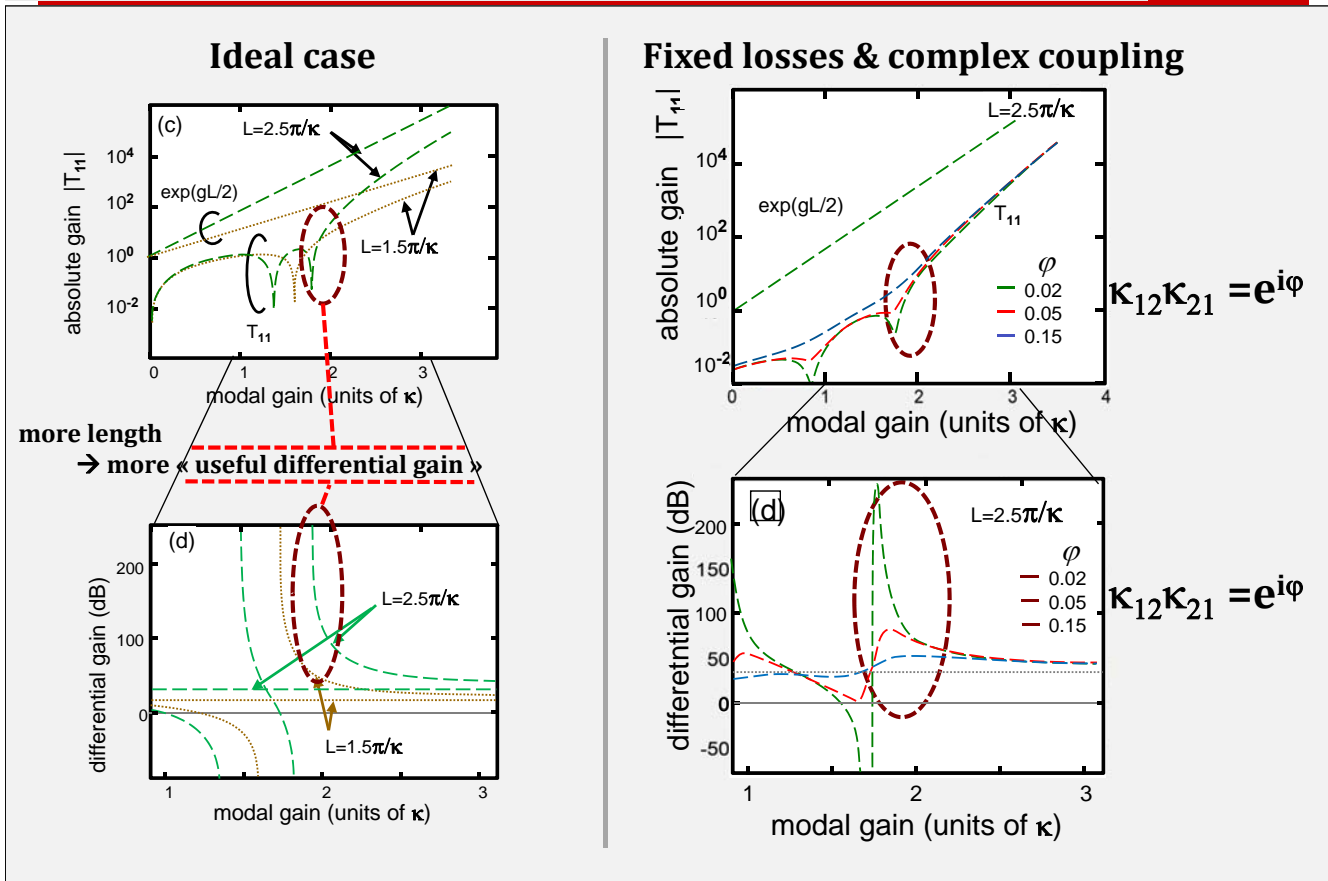
## Origin of complex coupling

microscopic picture



→ coupling value  $\kappa_{21}$  also depends on  $g_1$

# « Differential gain » : {Ideal} vs. {Fixed losses and Complex coupling}



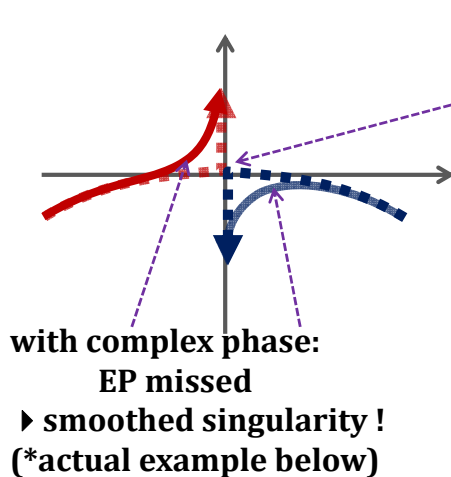
## PT-symmetry with plasmonics

► **Exceptional point smoothed when complex phase occurs in coupling**

$$\kappa_{12}\kappa_{21} = e^{i(\text{phase})}$$

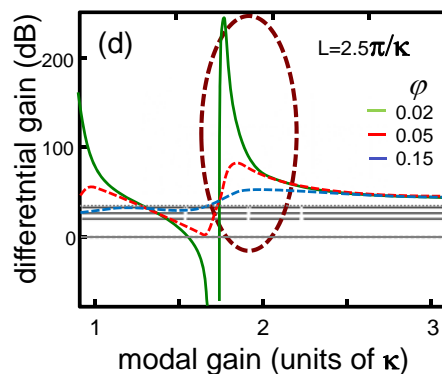
► **Small margin to retain high differential gain**

$$\kappa_{12}\kappa_{21} = e^{i(\text{small phase})}$$



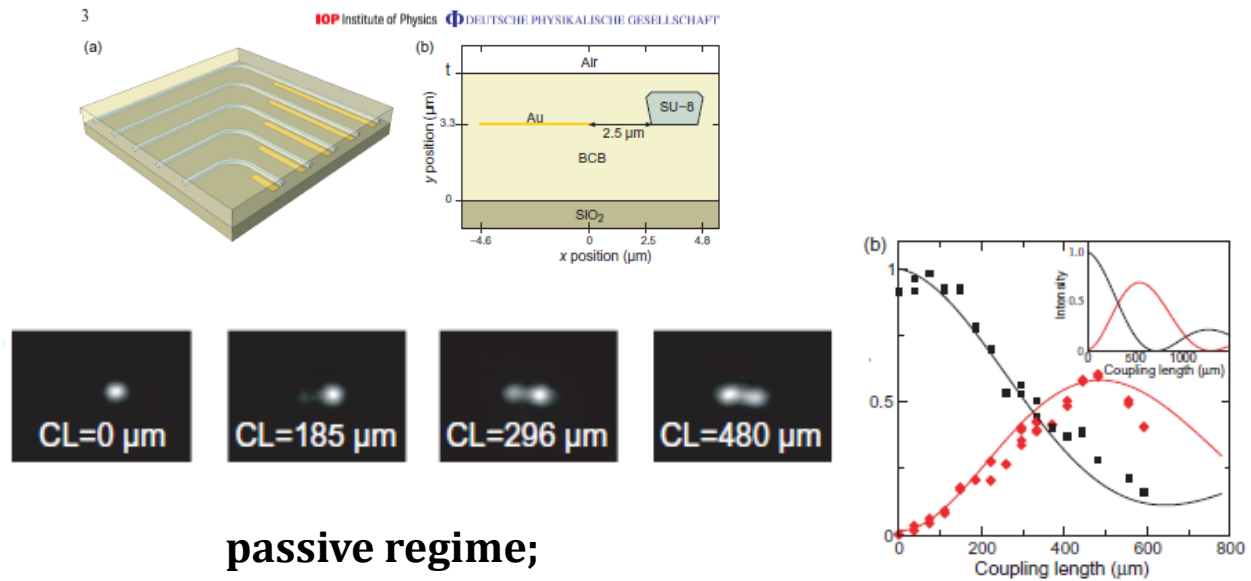
« Active-PT-structures »

**Ideal exceptional point**



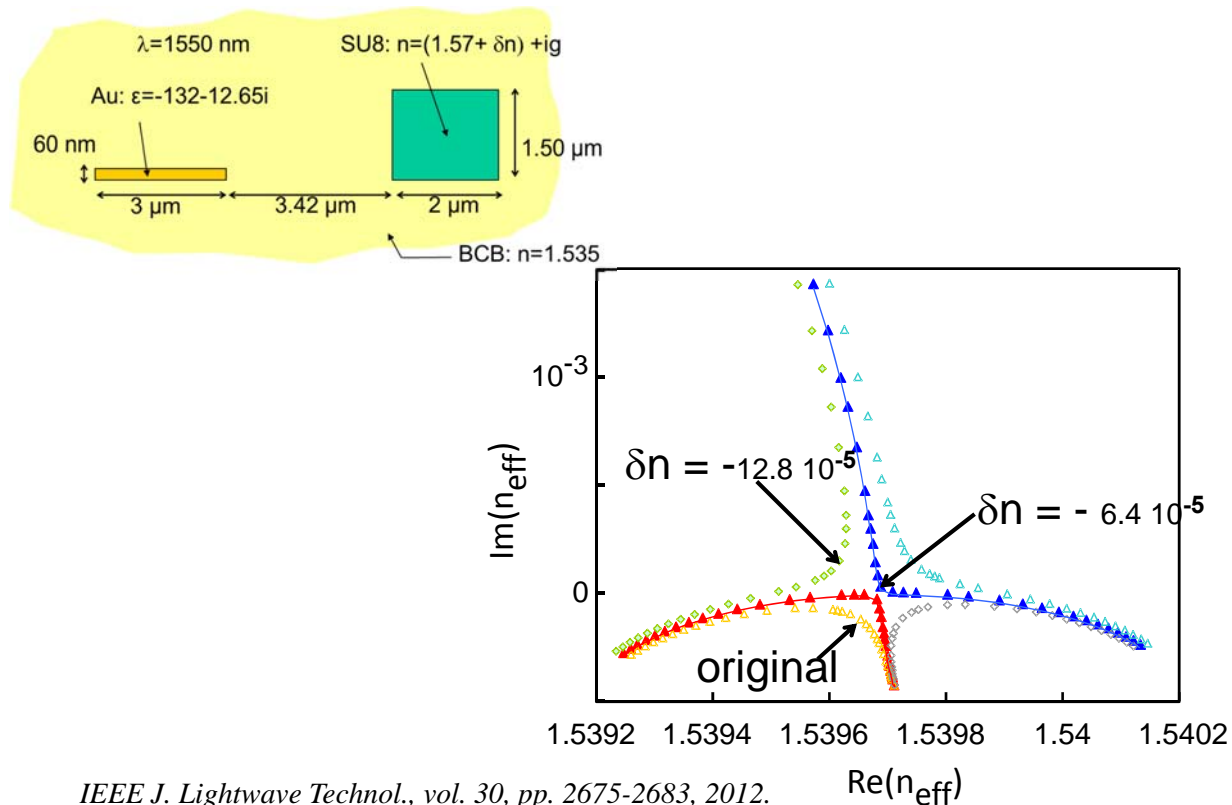
## polymer + Au-LRSPP

*New Journal of Physics* 11 (2009) 015002 A. Degiron et al. (@ Duke U.)



**passive regime;**  
**« Switching » based on**  
**detuning of  $\text{Re}(\epsilon)$**

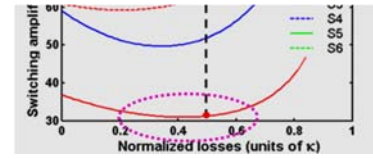
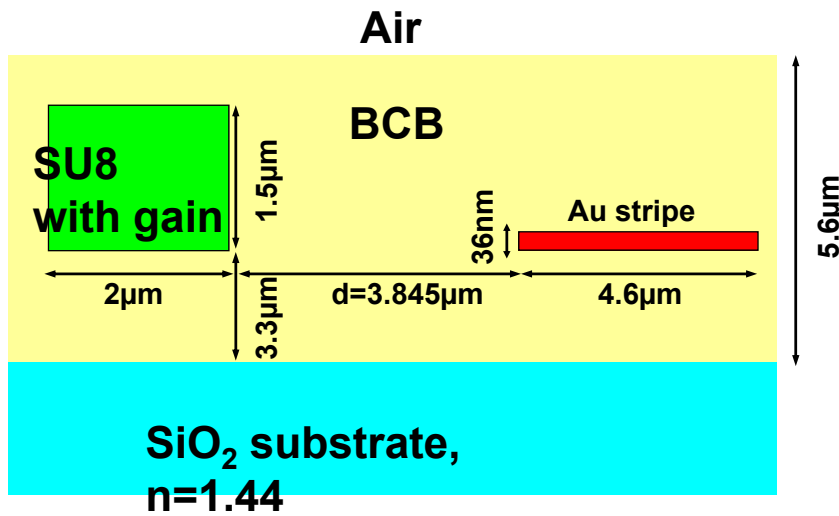
## « Healing » in this realistic architecture ?



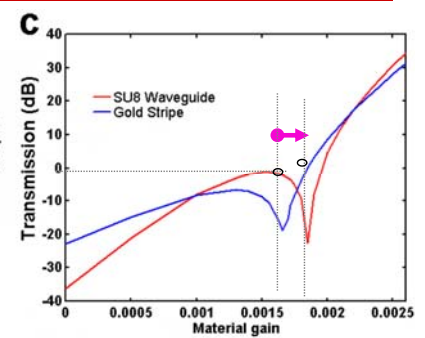
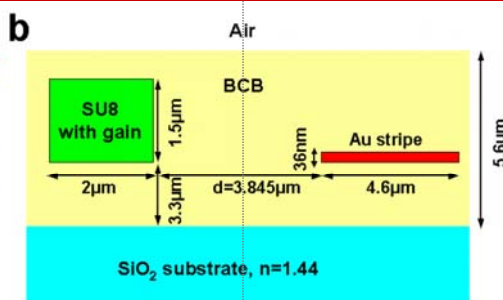
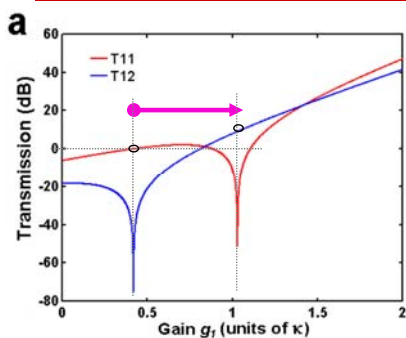
*IEEE J. Lightwave Technol.*, vol. 30, pp. 2675-2683, 2012.



Targetting  $\chi_2 = 0.42\kappa$ , loss/coupling ratio



## REALISTIC LRSP + SU8 SWITCH



### • CMT

- real  $\kappa$
- $\chi_2 = 0.42\kappa$ , fixed losses
- variable  $g_1$

### • FEM

- no « healing »

### real $\kappa$

- variable  $\text{Im}(\epsilon_{\text{SU8}})$

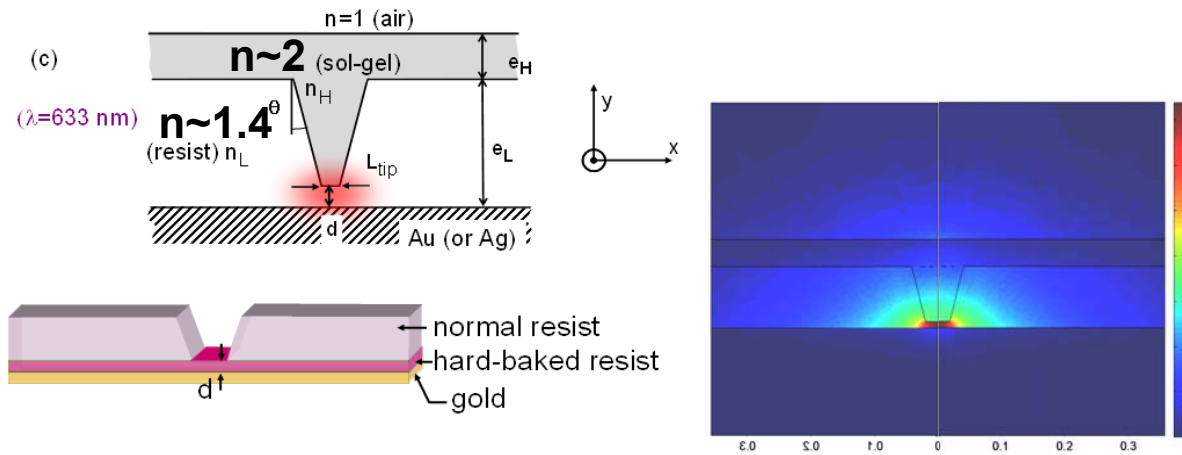
imperfections...

... less gain excursion

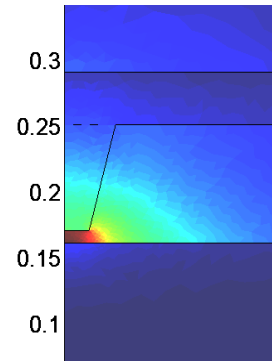
...

# Hybrid plasmonic waveguide : the «PIROW»

## PIROW: Plasmonic Inverse-Rib Optical Waveguide

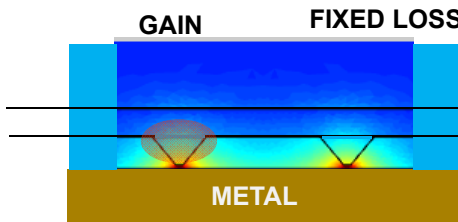


- JAP 2010, H. Benisty and M. Besbes  
E-field in 30-50 nm tip...
- Like Oulton's nanorod/spaser, but deterministic

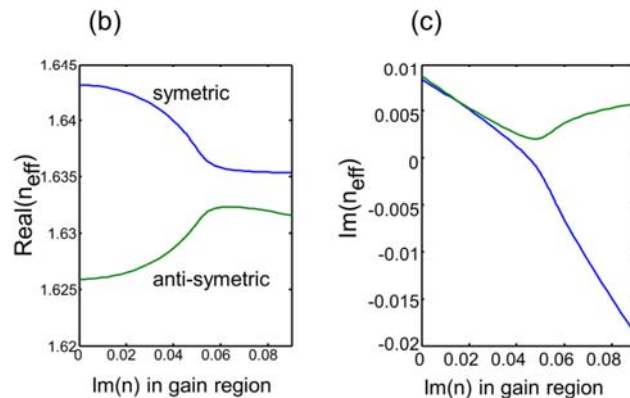
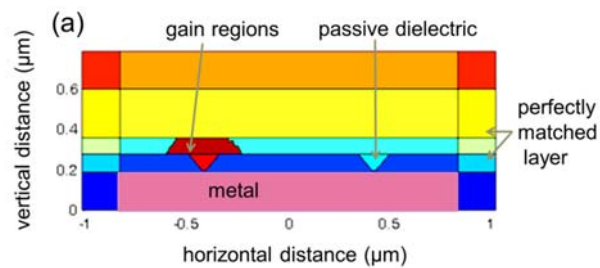


## Two coupled «PIROWs»

A normal one and a gain one

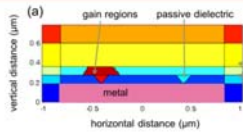


Can we have a good EP ?



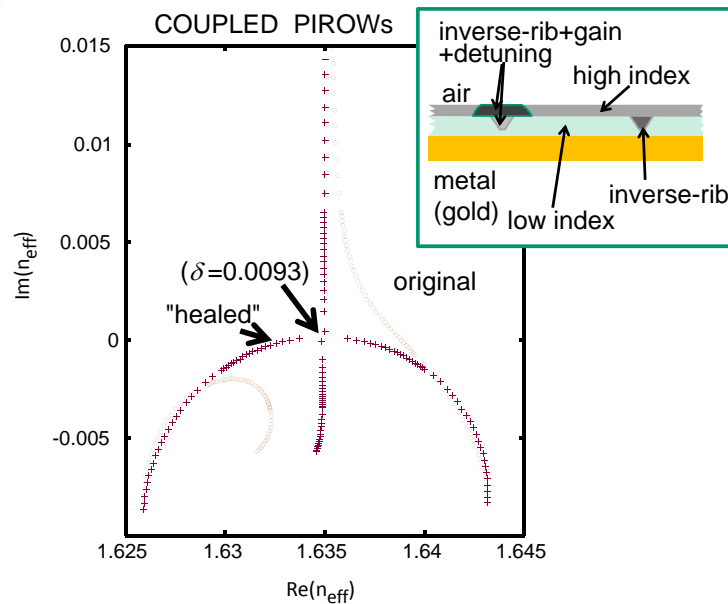
H. Benisty and M. Besbes,  
"Confinement and optical  
properties of the plasmonic  
inverse-rib waveguide,"  
*JOSA. B*, vol. 29,  
pp. 818-826, 2012.





$$(\tilde{n}_H) = (n_H) + ig(1 + i\alpha)$$

$$\alpha = -0.174$$



« Healing » obtained here by changing both  $\text{Im}(\epsilon)$  and  $\text{Re}(\epsilon)$ , the latter with a small factor... equivalent to detuning of waveguides with fixed  $\text{Re}(\epsilon)$  of EP...

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- Relevance for fast Quantum computation

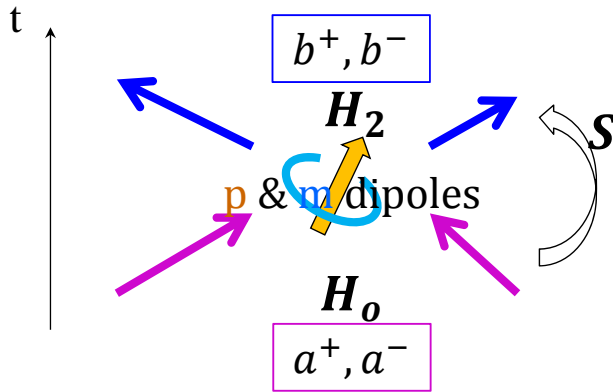
Jensen Li et al. work

$$\gamma_i = \gamma_i^s + \gamma_i^{\text{loss}}$$

Kang et al. (Jensen Li)  
PRA 87, 053824 (2013)

$$\mathbf{H}_n = \begin{pmatrix} \omega_1 - i\gamma_1 & \kappa \\ \kappa & \omega_2 - i\gamma_2 \end{pmatrix} + n \begin{pmatrix} i\gamma_1^s & 0 \\ 0 & i\gamma_2^s \end{pmatrix}, \quad (4)$$

$n=0,1,2$

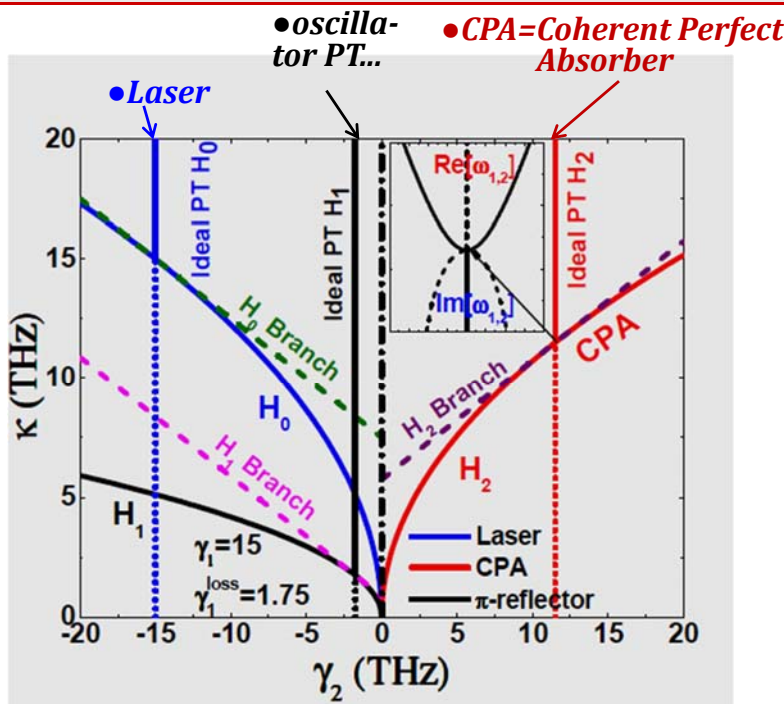


zero → CPA

$$\det(\mathbf{S}) = \frac{\det(\mathbf{H}_2 - \omega\mathbf{I})}{\det(\mathbf{H}_0 - \omega\mathbf{I})}$$

pole → laser

## Unification ? Each $\mathbf{H}_n$ has its PT-sym event



Kang et al. (Jensen Li)  
PRA 87, 053824 (2013)

Then,  $\det(\mathbf{H}_1 - \omega\mathbf{I}) = 0$  implies a complete reflection  $r = -1$  and  $t = 0$  for one side of incidence.

•Description based on a « rich dipole » (2 coupled degrees of freedom)

# The brachistochrone problem...



Johann Bernoulli

*Acta Eruditorum* 1696  $\frac{\sin \theta}{v} = \frac{1}{v} \frac{dx}{ds} = \frac{1}{v_m}$

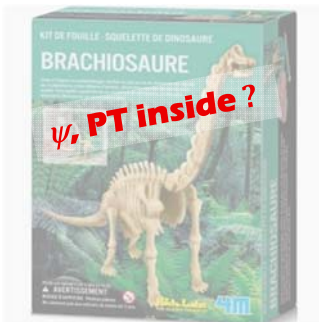


Quickest path  
with given  $g$  ?

~ "Quickest  
Hamiltonian"

[ Solution = Cycloid ]

# The quantum brachistochrone problem...



$$\psi_A \begin{pmatrix} 1 \\ 0 \end{pmatrix} \xrightarrow{H = \begin{bmatrix} s & r e^{-i\theta} \\ r e^{i\theta} & u \end{bmatrix}} \psi_B \begin{pmatrix} a \\ b \end{pmatrix}$$

$E_+ - E_- = \omega$   
«  $g$  »  $\leftrightarrow$  constraint on eigenvalues

[ Solution ~Rabi  $\pi/2$  oscillation]

$$\tau = \pi \hbar / \omega$$

**Faster than Hermitian Quantum Mechanics**

Bender et al...  
PRL 98, 040403 (2007)

**New inner product**

$$C^2 = 1, \quad [C, H] = 0, \quad \text{and} \quad [C, \mathcal{PT}] = 0.$$

$$C = \frac{1}{\cos\alpha} \begin{pmatrix} i \sin\alpha & 1 \\ 1 & -i \sin\alpha \end{pmatrix}$$

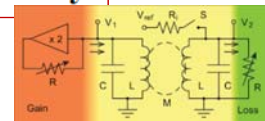
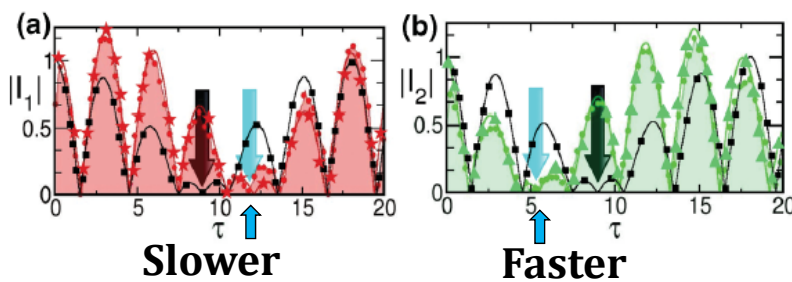
Seems to depend on H !

. If so, this would limit the applicability of a Hilbert-space worm-hole to improve quantum algorithms.

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T. Kottos, Wesleyan Univ....

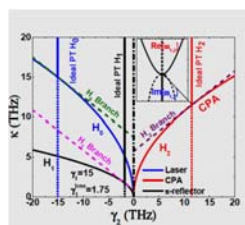
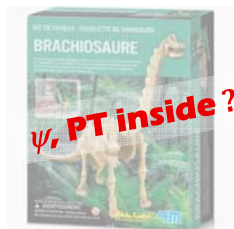
**Bypassing the bandwidth theorem with  $\mathcal{PT}$  symmetry**



$$\tau_{\text{fpt}} = \frac{1}{\delta\omega} \left[ \pi \pm \arccos \left( \frac{\delta\omega^2 - \gamma^2}{\delta\omega^2 + \gamma^2} \right) \right]$$

**CONCLUSION & PERSPECTIVES**

*Attractive concepts from classical to quantum*



*Quite some potential to unify several fields using Gain & "Phase-transitions" (lasers, CPA, EIT, metamaterial ?, strong coupling ?,...)*



*Can be combined (...with care...) with plasmonics to yield singularity from losses !*