Spectral embedding of graphs

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These lecture notes introduce the spectral embedding of graphs, where each node is represented by a vector of low dimension using the spectral decomposition of the Laplacian matrix. This embedding can in turn be used to apply classical learning techniques, either semi-supervised (with labels attached to some nodes) or unsupervised (like ranking or clustering nodes). We shall see that the spectral embedding of graphs is closely related to random walks in the graph. The analogy with various fields of physics, like thermodynamics, mechanics and electricity, will also emerge naturally.

These notes are mainly based on [1, 2, 3, 5]. We also refer the reader to [4] for an overview on spectral clustering.

1 Notion of embedding

Consider a weighted, undirected graph G = (V, E) of n nodes and m edges. Without loss of generality, we assume that $V = \{1, \ldots, n\}$. The weights are non-negative and correspond to the strengths of the links between nodes. The graph is assumed to be connected and without self-loops. We denote by A the weighted adjacency matrix of the graph, i.e., A_{ij} is the weight of the edge between nodes i and j, if any, and is equal to 0 otherwise. We denote by 1 the vector of ones and by w = A1 the vector of node weights (i.e., sums of the weights of incident edges); for unit edge weights, w is the vector of node degrees.

We aim at representing the graph in some Euclidian space of low dimension, say \mathbb{R}^k with $k \ll n$. Specifically, each node $i \in V$ is represented by some vector $x_i \in \mathbb{R}^k$. The structure of the graph must be encoded in its representation x_1, \ldots, x_n in the sense that two "close" nodes i, j in the graph should correspond to two "close" vectors x_i, x_j in the embedding space.

2 Random walk

Consider a random walk in the graph G with a probability of moving from node i to node j equal to A_{ij}/w_i . Let X_0, X_1, X_2, \ldots be the sequence of nodes visited by the random walk. This defines an irreducible Markov chain on $\{1, \ldots, n\}$ with transition matrix $P = D^{-1}A$, where D = diag(w). We have for all $t \ge 1$:

$$\forall i = 1, \dots, n, \quad P(X_t = i) = \sum_{j=1}^n P(X_{t-1} = j) P_{ji}.$$

Denoting the distribution of X_t as a vector π_t , we get:

$$\pi_t^T = \pi_{t-1}^T P,\tag{1}$$

so that

$$\pi_t^T = \pi_0^T P^t$$

where π_0 is the initial distribution. If the graph is strongly connected and aperiodic (that is, the largest common divisor of the cycle lengths is equal to 1), the following limit exists and is unique:

$$\pi = \lim_{t \to +\infty} \pi_t. \tag{2}$$

This is the stationary distribution, which satisfies the balance equations:

$$\pi^T = \pi^T P. \tag{3}$$

In particular, π is the unique left eigenvector of P for the eigenvalue 1 such that $\pi^T 1 = 1$ (observe that P1 = 1, that is, 1 is the corresponding right eigenvector). The vector π gives the frequency of visits of the random walk to each node. Since the graph is undirected, we have $\pi \propto w$, i.e., nodes are visited in proportion to their weights. That the stationary distribution π is unique, given by $\pi = w/|w|$ where $|w| = w^T 1$, follows from the fact that the eigenvalue 1 of P is simple, as proved in section 4.

Remark 1 It can be shown that the sequence π_t converges to π at an exponential rate equal to the modulus of the second largest eigenvalue of P.

Return time. Let $P_i = P(\cdot|X_0 = i)$ and E_i the corresponding expectation. We denote by $\sigma_i = E_i(\tau_i^+)$ the mean return time to node *i*, with $\tau_i^+ = \min\{t \ge 1 : X_t = i\}$. Since π_i is the frequency of visits to node *i*, we have

$$\sigma_i = \frac{1}{\pi_i}.\tag{4}$$

This will be proved in section 5.

Hitting time, commute time, escape probability. Let $H_{ij} = E_i(\tau_j)$ be the mean hitting time of node j from node i, with $\tau_i = \min\{t \ge 0 : X_t = i\}$. Observe that $H_{ij} = 0$ for j = i. We denote by $\rho_{ij} = H_{ij} + H_{ji}$ the mean commute time between nodes i and j. The escape probability from node i to node j is $e_{ij} = P_i(\tau_j < \tau_i^+)$, for any $i \ne j$. This is the probability of hitting node j before returning to node i.

Proposition 1 We have:

$$\rho_{ij} = \frac{1}{\pi_i e_{ij}}.$$

Proof. Let $\tau_{ij} = \min\{t > \tau_i : X_t = j\}$ be the hitting time of node j after having visited node i. We have:

$$\begin{aligned} \rho_{ij} &= \mathbf{E}_{i}(\tau_{ji}), \\ &= \mathbf{E}_{i}(\tau_{i}^{+}) + \mathbf{E}_{i}(\tau_{ji} - \tau_{i}^{+}), \\ &= \mathbf{E}_{i}(\tau_{i}^{+}) + \mathbf{E}_{i}((\tau_{ji} - \tau_{i}^{+})\mathbf{1}_{\{\tau_{ji} > \tau_{i}^{+}\}}), \\ &= \mathbf{E}_{i}(\tau_{i}^{+}) + \mathbf{P}_{i}(\tau_{ji} > \tau_{i}^{+})\mathbf{E}_{i}(\tau_{ji} - \tau_{i}^{+}|\tau_{ji} > \tau_{i}^{+}), \\ &= \mathbf{E}_{i}(\tau_{i}^{+}) + \mathbf{P}_{i}(\tau_{j} > \tau_{i}^{+})\mathbf{E}_{i}(\tau_{ji}), \\ &= \mathbf{E}_{i}(\tau_{i}^{+}) + (1 - e_{ij})\rho_{ij}. \end{aligned}$$

The result then follows from (4).

Since the commute time is symmetric, in the sense that $\rho_{ij} = \rho_{ji}$ for each $i \neq j$, it follows from Proposition 1 that $\pi_i e_{ij} = \pi_j e_{ji}$: the frequency of *direct* paths (without return) from *i* to *j* is equal to the frequency of *direct* paths from *j* to *i*. This is in fact a direct consequence of the *reversibility* of the Markov chain [1]

3 Laplacian matrix

Let D = diag(w). The Laplacian matrix is defined by

$$L = D - A.$$

This is the discrete version of the usual Laplace operator. The diffusion governed by the heat equation can be used in semi-supervised learning tasks to propagate labels.

Heat equation. Consider some strict subset S of $\{1, \ldots, n\}$ and assume that the temperature of each node $i \in S$ is set at some fixed value T_i . We are interested in the evolution of the temperatures of the other nodes. Heat exchanges occur through each edge of the graph proportionally to the temperature difference between the corresponding nodes, with a coefficient equal to the weight of the edge, thus interpreted as *thermal conductivity*. Then,

$$\forall i \notin S, \quad \frac{dT_i}{dt} = \sum_{j=1}^n A_{ij}(T_j - T_i),$$

that is

$$\forall i \not\in S, \quad \frac{dT_i}{dt} = -(LT)_i,$$

where T is the vector of temperatures. This is the heat equation in discrete space. At equilibrium, T satisfies Laplace's equation:

$$\forall i \notin S, \quad (LT)_i = 0, \tag{5}$$

We say that the vector T is harmonic. With the boundary conditions T_i for all $i \in S$, this defines a Dirichlet problem. Observing that $D^{-1}L = I - P$, Laplace's equation can be written equivalently

$$\forall i \notin S, \quad T_i = (PT)_i. \tag{6}$$

Proposition 2 (Uniqueness) There is at most one solution to the Dirichlet problem.

Proof. Since P is a stochastic matrix, it follows from (6) that the temperature of node i is the weighted average of the temperatures of its neighbors.

We first prove that the maximum and the minimum of the vector T are achieved on the boundary, that is for nodes in S. Let i be any node such that T_i is maximum. If $i \notin S$, it follows from (6) that T_j is maximum for all neighbors j of i. If no such node belongs to S, we apply again this argument until we reach a node in S. Such a node exists because the graph is connected. It achieves the maximum of the vector T. The proof is similar for the minimum.

Now consider two solutions T, T' to Laplace's equation. Then $\delta = T' - T$ is a solution of Laplace's equation with the boundary condition $\delta_i = 0$ for all $i \in S$. We deduce that $\delta_i = 0$ for all i (because both the maximum and the minimum are equal to 0), that is T' = T.

Now let $\tau_S = \min\{t \ge 0 : X_t \in S\}$ be the hitting time of the set S. Define:

$$P_{ij}^S = \mathbf{P}_i(\tau_j = \tau_S)$$

This is the probability that the random walker first hits S in node j when starting from node i. Observe that P^S is a stochastic matrix. In particular, $P_{ij}^S = \delta_{ij}$ (Kronecker delta) for all $i \in S$. By first-step analysis, we have:

$$\forall i \notin S, \quad P_{ij}^S = \sum_{k=1}^n P_{ik} P_{kj}^S. \tag{7}$$

Proposition 3 (Existence) The solution to the Dirichlet problem is

$$\forall i \notin S, \quad T_i = \sum_{j \in S} P_{ij}^S T_j.$$
(8)

Proof. The vector T defined by (8) satisfies:

$$\forall i \notin S, \quad \sum_{j=1}^{n} P_{ij}T_j = \sum_{j=1}^{n} P_{ij} \sum_{k \in S} P_{jk}^S T_k = \sum_{k \in S} P_{ik}^S T_k = T_i,$$

where we have used (7). Thus T satisfies (6). The proof then follows from Proposition 2.

4 Spectral analysis

The Laplacian matrix L is positive semi-definite:

Proposition 4 We have:

$$\forall v \in \mathbb{R}^n, \quad v^T L v = \sum_{i < j} A_{ij} (v_i - v_j)^2.$$

Proof. For all $v \in \mathbb{R}^n$,

$$v^{T}Lv = v^{T}(D - A)v,$$

= $\sum_{i,j=1}^{n} w_{i}v_{i}^{2} - \sum_{i,j=1}^{n} v_{j}A_{ij}v_{i},$
= $\sum_{i,j=1}^{n} A_{ij}v_{i}(v_{i} - v_{j}),$
= $\frac{1}{2}\sum_{i,j=1}^{n} A_{ij}(v_{i} - v_{j})^{2},$
= $\sum_{i < j} A_{ij}(v_{i} - v_{j})^{2}.$

The spectral theorem yields

$$=V\Lambda V^{T},$$
(9)

where $\Lambda = \text{diag}(\lambda_1, \ldots, \lambda_n)$ is the diagonal matrix of eigenvalues of L, with $0 \leq \lambda_1 \leq \lambda_2 \leq \ldots \leq \lambda_n$, and $V = (v_1, \ldots, v_n)$ is the matrix of corresponding eigenvectors, with $V^T V = I$. In view of Proposition 4, $v^T L v = 0$ implies $v \propto 1$ (recall that the graph is connected) so that $\lambda_1 = 0 < \lambda_2$ and $v_1 = 1/\sqrt{n}$. This proves in turn that the eigenvalue 1 of P is simple, since Pv = v if and only if Lv = 0.

L

A mechanical system. Consider n points of unit mass where points i and j are linked by a spring of stiffness A_{ij} following Hooke's law (i.e., force proportional to the distance). Now if the points are located according to some vector $v \in \mathbb{R}^n$ along a line, the potential energy accumulated in the springs is:

$$\frac{1}{2} \sum_{i < j} A_{ij} (v_i - v_j)^2,$$

that is $\frac{1}{2}v^T L v$ in view of Proposition 4.

We impose that the moment of inertia of the system (for a rotation around the origin) is equal to 1, that is $v^T v = 1$. Clearly, the vector v that minimizes the potential energy is $v = v_1$ (the corresponding potential energy is null). Now if we impose $1^T v = 0$, meaning that the centre of mass is at the origin, we obtain $v = v_2$ and $v^T L v = \lambda_2$, so that the eigenvalue λ_2 corresponds to twice the minimum value of potential energy. This is a consequence of the following characterization of the spectrum of the Laplacian.

Theorem 1 For all $k = 1, \ldots, n$,

$$\lambda_k = \min_{\substack{v:v^T v = 1 \\ v_1^T v = 0, \dots, v_{k-1}^T v = 0}} v^T L v,$$
(10)

the minimum being attained for $v = v_k$.

Proof. Let $v \in \mathbb{R}^n$ such that $v^T v = 1$. The vector $x = V^T v$, giving the coordinates $x_1 = v_1^T v, \ldots, x_n = v_n^T v$ of v in the basis of eigenvectors, satisfies:

$$x^T \Lambda x = v^T V \Lambda V^T v = v^T L v$$
 and $x^T x = v^T V V^T v = 1$,

so that the optimization problem (10) is equivalent to:

$$\min_{\substack{x:x^Tx=1\\x_1=0,\dots,x_{k-1}=0}} x^T \Lambda x.$$

The result then follows from the equality:

$$x^T \Lambda x = \sum_{i=1}^n \lambda_i x_i^2.$$

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Assume the system has a uniform circular motion around its center of mass, taken as the origin, so that $1^T v = 0$, with $v \neq 0$. Let ω be the angular velocity of the system. By Newton's second law of motion, the system is in equilibrium if and only if

$$\forall i = 1, \dots, n, \quad \sum_{j=1}^{n} A_{ij}(v_j - v_i) = -v_i \omega^2,$$

that is

$$Lv = \omega^2 v$$

This means that v is an eigenvector of L (different from v_1 since $1^T v = 0$) with eigenvalue ω^2 . In particular, the only possible values of angular velocity are $\sqrt{\lambda_2}, \ldots, \sqrt{\lambda_n}$. Moreover,

$$v^T L v = v^T v \omega^2,$$

where $v^T v$ is the moment of inertia of the system. For a unit moment of inertia $v^T v = 1$, we obtain:

$$v^T L v = \omega^2.$$

Thus the eigenvalues $\lambda_2, \ldots, \lambda_n$ are the squares of the possible values of angular velocities (the first eigenvalue $\lambda_1 = 0$ corresponding to the absence of rotation) and the eigenvectors v_2, \ldots, v_n are the corresponding equilibriums with unit moments of inertia.

5 Spectral embedding

Let $L^+ = V\Lambda^+ V^T$ be the pseudo-inverse of L, with $\Lambda^+ = \operatorname{diag}\left(0, \frac{1}{\lambda_2}, \dots, \frac{1}{\lambda_n}\right)$.

Proposition 5 We have:

$$LL^{+} = L^{+}L = I - \frac{11^{T}}{n}.$$

Proof. The proof follows from the fact that $v_1 = 1/\sqrt{n}$ on observing that

$$LL^{+} = L^{+}L = V\Lambda^{+}\Lambda V^{T} = \sum_{k=2}^{n} v_{k}v_{k}^{T} = I - v_{1}v_{1}^{T}.$$

Let:

$$X = \sqrt{|w|}Z(I - \pi 1^T),$$

where

$$Z = \sqrt{\Lambda^+} V^T.$$

Consider the embedding $X = (x_1, \ldots, x_n)$ of the graph, where node *i* is represented by the vector $x_i \in \mathbb{R}^n$. Observe that the first row of X is null so that only n-1 coordinates are informative. The embedding X is a shifted, rescaled version of Z so that:

$$X\pi = 0.$$

The Gram matrix of Z is the pseudo-inverse of the Laplacian L:

$$Z^T Z = V \Lambda^+ V^T = L^+$$

We deduce the Gram matrix of X,

$$G = X^{T}X = |w|(I - 1\pi^{T})L^{+}(I - \pi 1^{T}).$$
(11)

Observe that

$$G\pi = 0. \tag{12}$$

For any matrix M, we denote by d(M) the diagonal matrix with the same diagonal as that of M.

Random walk. By one-step analysis, the mean hitting time of node j from node i satisfies:

$$H_{ij} = \begin{cases} 0 & \text{if } i = j, \\ 1 + \sum_{k=1}^{n} P_{ik} H_{kj} & \text{otherwise.} \end{cases}$$
(13)

We deduce that the matrix $(I - P)H - 11^T$ is diagonal. Equivalently, the matrix $LH - w1^T$ is diagonal.

Lemma 1 There is at most one matrix H such that d(H) = 0 and the matrix $LH - w1^T$ is diagonal.

Proof. Let H, H' be two such matrices and $\Delta = H - H'$. We have $L\Delta = 0$ so that each column of Δ is either null or proportional to 1. Since $d(\Delta) = 0$, we get $\Delta = 0$, that is H' = H.

Theorem 2 We have:

$$H = 11^T d(G) - G, (14)$$

where $G = X^T X$ is the Gram matrix of X.

Proof. Using the fact that L1 = 0, the matrix H defined by (14) satisfies:

$$LH = -LG,$$

$$= -|w|L(I - 1\pi^{T})L^{+}(I - \pi 1^{T}),$$

$$= -|w|LL^{+}(I - \pi 1^{T}),$$

$$= -|w|(I - \frac{11^{T}}{n})(I - \pi 1^{T}),$$

$$= -|w|(I - \pi 1^{T}),$$

$$= -|w|I + w1^{T},$$
(15)

so that the matrix $LH - w1^T$ is diagonal. Since d(H) = 0, the proof follows from Lemma 1.

Observe that the mean return time to node i satisfies:

$$\sigma_i = 1 + \sum_{j=1}^n P_{ij} H_{ji},$$

so that the corresponding vector σ forms the diagonal of the matrix $PH + 11^T$. The following result then proves (4).

Proposition 6 We have:

$$d(PH+11^T) = \operatorname{diag}(\pi)^{-1}.$$

Proof. In view of (15), $(I - P)H = -\text{diag}(\pi)^{-1} + 11^T$, and the result follows on observing that d(H) = 0. \Box

Let h_i be the mean hitting time of node *i* in stationary regime (that is, starting from a node chosen at random from the stationary distribution):

$$h_i = \sum_{j=1}^n \pi_j h_{ji}.$$

The corresponding vector h satisfies $h^T = \pi^T H$. In view of Theorem 2 and equation (12),

$$h^T = 1^T d(G),$$

In particular,

 $H = 1h^T - G,$

that is

$$x_i^T x_j = h_j - H_{ij} = h_i - H_{ji},$$

and, since $h_{ii} = 0$,

$$||x_i||^2 = h_i$$

In particular, the mean commute time between nodes i and j is given by:

$$\rho_{ij} = H_{ij} + H_{ji} = ||x_i - x_j||^2.$$

6 Electric networks

Consider the electric network induced by the graph, with a resistor of conductance A_{ij} between nodes *i* and *j*. We denote by 1_i the unit vector on component *i*.

Effective conductance, effective resistance. For any distinct nodes s, t, assume the electric potentials of s and t are set to 1 and 0, respectively. Let U_i be the electric potential of any node i. We have $U_s = 1$ and $U_t = 0$. By Ohm's law, the current that flows from i to j is

$$A_{ij}(U_i - U_j).$$

By Kirchoff's law, the net current at any node $i \neq s, t$ is null. We get

$$\sum_{j=1}^{n} A_{ij}(U_i - U_j) = 0,$$

that is $(LU)_i = 0$. Thus the vector of electric potentials U is harmonic. Moreover, $(LU)_s + (LU)_t = 0$, so that $LU = \alpha(1_s - 1_t)$ for some constant α , equal to the current flowing from s to t.

Proposition 7 We have:

$$U_i = \frac{(x_i - x_t)^T (x_s - x_t)}{||x_s - x_t||^2}.$$

Proof. In view of Proposition 5,

$$\left(I - \frac{11^T}{n}\right)U = L^+ L U = \alpha L^+ (1_s - 1_t),$$

that is

$$U = \alpha L^+ (1_s - 1_t) + \beta 1,$$

with $\beta = 1^T U/n$. We obtain:

$$\alpha z_s^T (z_s - z_t) + \beta = 1,$$

$$\alpha z_t^T (z_s - z_t) + \beta = 0,$$

so that

$$\alpha = \frac{1}{||z_s - z_t||^2}, \quad \beta = -\frac{z_t^T(z_s - z_t)}{||z_s - z_t||^2}$$

Finally,

$$U_i = \frac{(z_i - z_t)^T (z_s - z_t)}{||z_s - z_t||^2},$$

and the proof follows from the fact that $x_i - x_j = \sqrt{|w|}(z_i - z_j)$ for all i, j = 1, ..., n.

The current α flowing from s to t, which is the current induced by a unit electric potential, is called the *effective conductance* between s and t. In view of the proof of Proposition 7, we have:

$$\alpha = \frac{1}{||z_s - z_t||^2} = \frac{|w|}{\rho_{st}},$$

so that the effective conductance between s and t is proportional to $1/\rho_{st}$, the inverse of the mean commute time of the random walk between s and t. Equivalently, the mean commute time ρ_{st} between nodes s and t can be interpreted as the *effective resistance* between s and t in the electric network, in some arbitrary unit.

Thompson's principle. The energy dissipation through any transistor is the product of voltage and current (both in absolute value), that is

$$A_{ij}(U_j - U_i)^2$$

between nodes i and j. We obtain the total energy dissipation:

$$E = \frac{1}{2} \sum_{i,j=1}^{n} A_{ij} (U_j - U_i)^2,$$

that is, in view of Proposition 4,

$$E = U^T L U.$$

Thompson's principle states that the potential vector U minimizes energy dissipation. Taking the derivative in U_i , we obtain:

$$\sum_{j=1}^{n} A_{ij}(U_j - U_i) = 0,$$

that is $(LU)_i = 0$, which is Laplace's equation.

Interpretation of voltage and current. Observe that the electric potential is the solution to the heat equation with $T_s = 1$ and $T_t = 0$. In view of (8), we have $U_i = P_{is}^S$, i.e., the electric potential of any node is the probability that the random walk reaches node s before node t. Thus everything happens as if each electron were a random walker in the graph. We shall see that the current between two nodes can be interpreted as the net flow of electrons between these two nodes.

For convenience, we consider positive particles starting from node s and captured by node t (thus in the direction of the current) instead of electrons starting from node t and captured by node s, but the interpretation is exactly the same. Consider the path of a particule starting from node s before it is captured by node t. Let N_i be the mean number of times it visits node i before being captured by node t. We take the initial state into account in the number of visits so that, by Proposition 1,

$$N_s = \frac{1}{e_{st}} = \pi_s \rho_{st},$$

while $N_t = 0$. For any $i \neq s, t$, we have by one-step analysis,

$$N_i = \sum_{j=1}^n P_{ji} N_j.$$

Using the local balance equation $\pi_i P_{ij} = \pi_j P_{ji}$, we get

$$\frac{N_i}{\pi_i} = \sum_{j=1}^n P_{ij} \frac{N_j}{\pi_j}.$$

We deduce that the vector U defined by

$$U_i = \frac{N_i}{\pi_i \rho_{st}}$$

is harmonic, with $U_s = 1$ and $U_t = 0$. This is the electric potential. The net current from node i to node j is

$$A_{ij}(U_i - U_j) = \frac{1}{\rho_{st}} \left(\frac{N_i}{\pi_i} A_{ij} - \frac{N_j}{\pi_j} A_{ji} \right) = \frac{|w|}{\rho_{st}} (N_i P_{ij} - N_j P_{ji})$$

This is the net frequency of particle moving from node i to node j, with a flow of particles entering the network at node s at rate

$$\alpha = \frac{|w|}{\rho_{st}},$$

which is the current flowing from node s to node t.

General solution. Now consider the case where the electric potential of node s is set to 1 while those of K other nodes, say t_1, \ldots, t_K , are set to 0. The following result extends Proposition 7.

Proposition 8 We have:

$$U_{i} = \sum_{k=1}^{K} \alpha_{k} (x_{i} - x_{t_{l}})^{T} (x_{s} - x_{t_{k}}),$$

where l is an arbitrary element of $\{1, \ldots, k\}$ and the vector $\alpha = (\alpha_1, \ldots, \alpha_K)^T$ is the unique solution to the equation $M\alpha = |w|1$, with M the Gram matrix of the vectors $(x_s - x_{t_1}, \ldots, x_s - x_{t_K})$.

Proof. Let α_k be the current going out of node t_k , for $k = 1, \ldots, K$. Then $\sum_{k=1}^{K} \alpha_k$ is the current entering node s and

$$LU = \sum_{k=1}^{K} \alpha_k (1_s - 1_{t_k}).$$

By Proposition 5,

$$\left(I - \frac{11^T}{n}\right)U = \sum_{k=1}^K \alpha_k L^+ (1_s - 1_{t_k}),$$

that is

$$U = \sum_{k=1}^{K} \alpha_k L^+ (1_s - 1_{t_k}) + \beta 1,$$

with $\beta = 1^T U/n$. We get:

$$\sum_{k=1}^{K} \alpha_k z_s^T (z_s - z_{t_k}) + \beta = 1,$$

$$\sum_{k=1}^{K} \alpha_k z_{t_l}^T (z_s - z_{t_k}) + \beta = 0, \quad l = 1, \dots, K.$$

In particular,

$$\sum_{k=1}^{K} \alpha_k (z_s - z_{t_l})^T (z_s - z_{t_k}) = 1, \quad l = 1, \dots, K,$$

so that α is the unique solution to the equation $M\alpha = |w|1$ (recall that $x_j - x_i = \sqrt{|w|}(z_j - z_i)$ for all i, j). The result then follows easily.

Similarly, the electric potential U is the solution to the heat equation with $U_s = 1$ and $U_{t_1}, \ldots, U_{t_K} = 0$. It follows from (8) that $U_i = P_{is}^S$, the probability that a random walk starting from i hits the set $S = \{s, t_1, \ldots, t_K\}$ in s. Thus applying Proposition 8 to each $s \in S$ provides the full matrix P^S and thus the solution for any boundary condition. Specifically, setting the electric potential U_i of node i, for each $i \in S$, yields the solution:

$$\forall i \notin S, \quad U_i = \sum_{j \in S} P_{ij}^S U_j.$$

7 Applications

Finally, we show how to apply previous results to problems of node ranking and clustering. The first step consists in computing the embedding of the graph, $X = (x_1, \ldots, x_n)$:

- Embedding -

Parameter: k, dimension of the embedding

- 1. Check that the graph is connected
- 2. Form the Laplacian L = D A
- 3. Compute v_1, \ldots, v_k , the k eigenvectors of L associated with the lowest eigenvalues, $\lambda_1 \leq \ldots \leq \lambda_k$
- 4. Compute $Z = \text{diag}\left(\frac{1}{\sqrt{\lambda_2}}, \dots, \frac{1}{\sqrt{\lambda_k}}\right) (v_2, \dots, v_k)^T$

5. Return
$$X = \sqrt{|w|Z(I - \pi 1^T)}$$
 where $w = w1$ and $\pi = w/|w|$

Ranking. A first way to rank nodes is to consider their *centrality*, in terms of mean hitting time: the more central the node, the shorter time it takes on average for a random walk to hit this node. By the results of section 5, we get the following ranking, the most central nodes appearing first:

- Centrality -

Output: nodes in increasing order of $||x_i||^2$

In practice, it is often interesting to rank nodes relative to another node. We then rank nodes with respect to their *local centrality*, defined as the mean hitting time from the node of interest. This approach easily extends to a set of nodes. By the results of section 5, we get:

Local centrality Input: s, node of interest Output: nodes in increasing order of $x_i^T(x_i - x_s)$

It may also be interesting to include, in addition to the node of interest, a repulsive node. We can then rank nodes with respect to their *directional centrality*, corresponding to the probability to hit the node of interest before the repulsive node (which can be interpreted as an electric potential in view of the results of section 6). Again, the approach easily extends to a set of repulsive nodes.

Directional centrality —

Input: *s*, node of interest; *t*, repulsive node **Output:** nodes in increasing order of $x_i^T(x_t - x_s)$

Clustering. For clustering the nodes of the graph, one may apply classical clustering techniques in Euclidian spaces. The most popular algorithm is K-means (or variants like fuzzy K-means to enable overlaps), which consists in finding the partition C_1, \ldots, C_K of $\{1, \ldots, n\}$ that minimizes:

$$J = \sum_{k=1}^{K} \sum_{i \in C_k} ||x_i - g(C_k)||^2,$$

where

$$g(C_k) = \frac{1}{|C_k|} \sum_{i \in C_k} x_i$$

is the centroid of cluster C_k . We have the following classical result:

Proposition 9 Let g be the centroid of n vectors x_1, \ldots, x_n . Then,

$$\sum_{i=1}^{n} ||x_i - g||^2 = \frac{1}{2n} \sum_{i,j=1}^{n} ||x_i - x_j||^2.$$

Proof. We have

$$\sum_{i=1}^{n} ||x_i - g||^2 = \sum_{i=1}^{n} x_i^T (x_i - g) = \sum_{i=1}^{n} ||x_i||^2 - \frac{1}{n} \sum_{i,j=1}^{n} x_i^T x_j = \frac{1}{2n} \sum_{i,j=1}^{n} ||x_i - x_j||^2.$$

In view of Proposition 9, we have

$$J = \sum_{k=1}^{K} \frac{1}{2|C_k|} \sum_{i,j \in C_k} ||x_i - x_j||^2.$$

Thus the cost function J can be interpreted, up to a factor n/2, as the mean square distance of a random point to another random point of the same cluster. In view of the results of section 5, the best clustering for the cost function J is that minimizing the mean commute time of the random walk between a random node and another node taken uniformly at random in the same cluster.

Hierarchical clustering. The main drawback of K-means is the need to specify the number of clusters K. An alternative consists in representing the graph by a binary tree, revealing the multi-scale structure of the graph, at different resolutions. This is known as hierarchical clustering. The most popular algorithm is the Ward method [6]. This is an agglomerative algorithm based on the sum of square errors, as in K-means. Appropriate cuts of the tree provide relevant clusterings of the graph. All these clusterings are encoded in a single data structure, a binary tree, with requires O(n) memory only.

References

- P. Brémaud. Markov chains: Gibbs fields, Monte Carlo simulation, and queues. Springer Science & Business Media, 2013.
- [2] F. R. Chung. Spectral graph theory. American Mathematical Soc., 1997.
- [3] L. Lovász. Random walks on graphs. Combinatorics, Paul Erdos is eighty, 1993.
- [4] U. Luxburg. A tutorial on spectral clustering. Statistics and Computing, 2007.
- [5] P. Snell and P. Doyle. Random walks and electric networks. Free Software Foundation, 2000.
- [6] J. H. Ward. Hierarchical grouping to optimize an objective function. Journal of the American statistical association, 1963.